

it has been shown that reaction 1 occurs at thermal collision energies. Currently available thermodynamic data are not of sufficient accuracy to indicate whether reaction 1 is endothermic or exothermic. However, it is clear from the results presented that this reaction must be exothermic to account for the observation of  $\text{CeO}_2^+$  at thermal collision energies and the observation of chemielectrons with energies of  $\sim 1$  eV. There is therefore clearly a need for more accurate thermodynamic data on the cerium-oxygen system. Nevertheless, the qualitative discussion presented in this paper demonstrates how chemielectron spectroscopy is capable of probing the potential energy surfaces of reacting systems and providing complementary information on molecular ions to that obtained from photoelectron spectroscopy.

Related work is currently in progress on other lanthanide plus oxygen chemiionization reactions. In particular, the  $\text{La} + \text{O}_2$ ,  $\text{Pr} + \text{O}_2$ , and  $\text{Nd} + \text{O}_2$  reactions are being studied, and for these reactions available cross sections<sup>12</sup> suggest that the analogue of reaction 6 is faster than in the  $\text{Ce} + \text{O}_2$  case.

**Acknowledgment.** We thank the SERC for financial assistance, and A.M.E. thanks the CEGB for the award of a CASE studentship. This work was also supported in part by the Air Force Office of Scientific Research (Grant No. AFOSR-89-0351) through the European Office of Aerospace Research (EOARD), United States Air Force.

Registry No. Ce, 7440-45-1;  $\text{O}_2$ , 7782-44-7.

## Hydrogen-Abstraction Reactivity of Excited-State Para-Substituted Benzyl Radicals in Solution at Room Temperature

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Received August 12, 1988

**Abstract:** In contrast to almost no H-abstraction reactivity of the ground-state para-substituted benzyl radicals ( $p$ -X-benzyls; X = CN, Cl, F, and  $\text{OCH}_3$ ) toward 1,4-cyclohexadiene (CHD), a significant fluorescence quenching by CHD in hexane at room temperature was confirmed for their fluorescent excited states, which may be assigned to  $2\text{B}_2$  for  $p$ -cyanobenzyl and  $1\text{A}_2$  for the other three  $p$ -X-benzyls. The quenching reveals a marked substituent effect: the quenching rate constant of  $3.2 \times 10^9 \text{ M}^{-1} \text{ s}^{-1}$  for  $p$ -cyanobenzyl fluorescence is much larger than that of  $9.4 \times 10^5 \text{ M}^{-1} \text{ s}^{-1}$  for  $p$ -methoxybenzyl fluorescence. The excited-state H-abstraction of  $p$ -cyanobenzyl from a H-donating CHD was directly demonstrated by the transient absorption induced upon the 308-nm pulse excitation of the ground-state  $p$ -cyanobenzyl. The marked substituent effect upon the quenching by CHD implies that  $1\text{A}_2$  and  $2\text{B}_2$  states have considerably different unpaired electron densities at the benzyl position.

Various multiphoton processes are operative in the organic photochemical reactions induced by intense laser pulses. One of the most efficient biphotonic processes is the photolysis of the transient species such as excited states and/or various reactive intermediates. Many free radicals are known as the reactive intermediate in the organic photochemistry. The study of the reactivity of the excited-state aromatic free radical would be thus expected to elucidate the laser-induced organic photochemistry, which is usually performed in the liquid phase at room temperature. Such a study (one of the photochemistry of transient species) is an attractive and current subject in the physical organic chemistry. The unimolecular rearrangement in the lowest excited doublet state has been confirmed for the very long-lived triphenylmethyl<sup>1</sup> and the stable perchlorotriphenylmethyl<sup>2</sup> radicals in solution at room temperature. Upon photoexcitation of unstable aromatic radicals (benzophenone ketyl<sup>3</sup> and the several arylmethyl radicals such as diphenylmethyl<sup>4</sup> and 1-naphthylmethyl<sup>5</sup>), furthermore, the intermolecular reactions with various substrates have been demonstrated to take place in solution at room temperature.<sup>3-7</sup> Long-lived excited-state radicals are generally preferable

for the intermolecular reactions in fluid solution.

It has been demonstrated that the fluorescent-state benzyl radical (the prototype of arylmethyl radicals) in solution at room temperature is very short-lived ( $\leq 1$  ns) owing to an extensively temperature-dependent nonradiative relaxation.<sup>8,9</sup> However, the fluorescence lifetimes of  $p$ -chlorobenzyl and  $p$ -methoxybenzyl in hexane at room temperature have recently been determined to be  $81 \pm 5$  ns and  $120 \pm 5$  ns,<sup>10</sup> which are much longer than that of benzyl. Such relatively long fluorescence lifetimes stimulated us to study the intermolecular reactivity of some para-substituted benzyls in the fluorescent excited state.

Based upon the kinetics and spectra of the fluorescence and absorption induced upon the 308-nm pulse excitation of the ground-state  $p$ -X-benzyl radical (X =  $\text{OCH}_3$ , F, Cl, or CN) in hexane at room temperature, we wish to report the quenching reaction of the excited-state  $p$ -X-benzyl by 1,4-cyclohexadiene known as a potential H donor. A pronounced substituent effect upon the quenching is discussed in terms of the  $1\text{A}_2 \leftrightarrow 2\text{B}_2$  alternation of the lowest excited doublet state.

### Experimental Section

$p$ -(Chloromethyl)anisole (4-methoxybenzyl chloride, Tokyo Kasei), 4-fluorobenzyl bromide (Aldrich), and 4-chlorobenzyl chloride (Nakarai) were used after distillation under reduced pressure.  $\alpha$ -Bromo- $p$ -tolunitrile (4-cyanobenzyl bromide, Aldrich) was purified by recrystallization from

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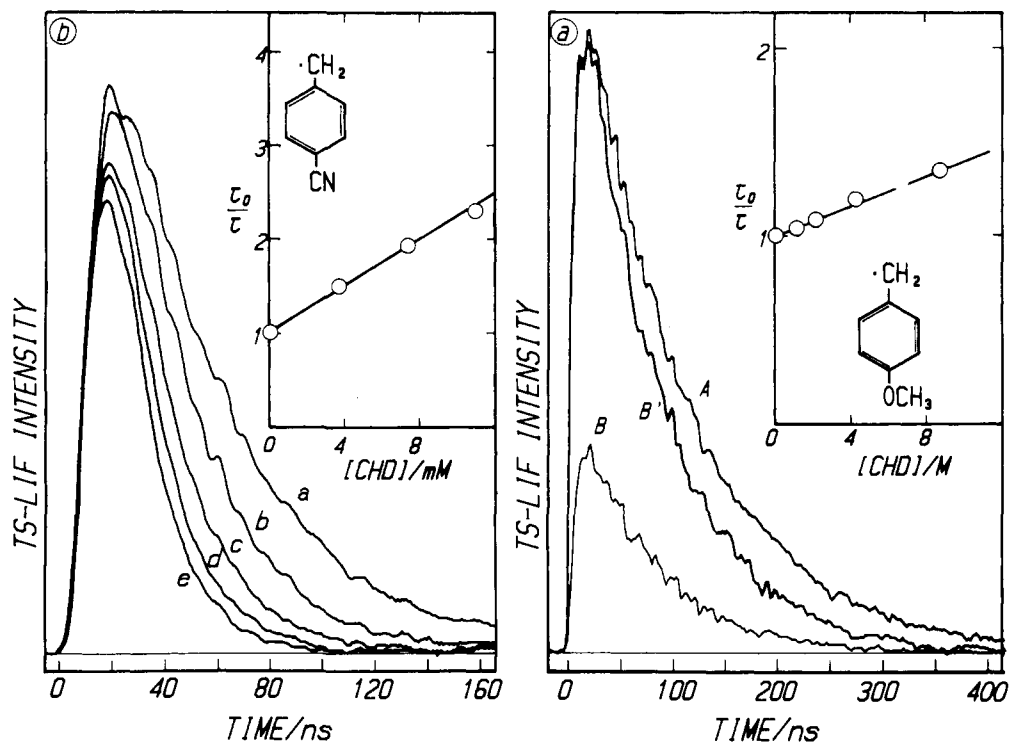
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**Figure 1.** Oscilloscope traces for the  $D_1 \rightarrow D_0$  fluorescence of *p*-methoxybenzyl (a) and *p*-cyanobenzyl (b) in the presence of CHD. The concentrations of CHD are 0 (a and A), 3.63 (b), 7.27 (c), 10.9 (d), 14.5 (e), and 4360 mM (B and B'). The intensities of fluorescence signals (A and B') are nearly normalized at maximum. The Stern-Volmer plots ( $\tau_0/\tau$  vs CHD concentration) are shown in the respective insets.

ethanol. Spectral grade hexane (Nakarai) without further purification was used as solvent to dissolve benzyl halides. 1,4-Cyclohexadiene (CHD) was purified by a silica column chromatography. Sample solutions were degassed by repeated freeze-pump-thaw cycles.

The 248-nm KrF laser pulse (14-ns fwhm, 75 mJ) was employed as the photolysis pulse to induce the homolytic scission of the  $\text{CH}_2\text{-X}$  bond in benzyl halide to yield a benzyl radical and a halogen atom, and the 308-nm XeCl laser pulse (13-ns fwhm, 80 mJ) was selected as the probe pulse to excite the radical in the ground state. Transient absorption and fluorescence spectra were measured by the laser-flash apparatus consisting of excimer lasers (Lambda Physik EMG 50E and 53MSC) with and without the monitoring pulsed xenon flash, respectively. The monitoring beam was cut off by an electromechanical shutter operating synchronously with the lasers. The delay time between the 248-nm and 308-nm pulses was controlled by a digital delay generator (BNC 7010) with the minimum resolution time of 100 ns. The electric signals from the photomultiplier (Hamamatsu Photonics R666 or R928) were digitized by a storage scope (Iwatsu TS 8123), and the kinetic and spectral analyses of the digitized data were carried out by a microcomputer (Fujitsu Micro 11BS or 16 $\beta$ ).

## Results and Discussion

High concentrations of free radical are often obtainable by the laser photolysis of its precursor in solution. As is distinct from the effective H-abstraction of alkyl radicals from H-containing solvents, most aromatic free radicals with unpaired electron delocalized over aromatic ring undergo various coupling reactions,<sup>11-15</sup> the rates of which are sufficiently slow to undertake the second pulse excitation. The  $D_1 \rightarrow D_0$  fluorescence and the  $D_n \leftarrow D_1$  absorption spectra were induced by the 308-nm XeCl laser pulse excitation of the transient para-substituted benzylic (*p*-X-benzyl) formed in the 248-nm KrF laser photolysis of *p*-X-benzyl halides in hexane at room temperature.<sup>9,10</sup> The fluorescence lifetimes ( $\tau_0$ ) were determined to be  $120 \pm 8$  ns (*p*-methoxybenzyl),  $200 \pm 8$  ns (*p*-fluorobenzyl),  $81 \pm 5$  ns (*p*-chlorobenzyl),

and  $58 \pm 4$  ns (*p*-cyanobenzyl), respectively.

**Fluorescence Quenching of Para-Substituted Benzyl Radicals by 1,4-Cyclohexadiene.** The kinetic electron paramagnetic resonance experiment has demonstrated<sup>16</sup> that the ground-state benzyl radical does not abstract hydrogen from 1,4-cyclohexadiene (CHD), implying that the rate constant is less than  $10^2 \text{ M}^{-1} \text{ s}^{-1}$  at 300 K, in contrast to the appreciable H-abstraction of alkyl radicals from CHD with rate constants of  $10^4\text{--}10^5 \text{ M}^{-1} \text{ s}^{-1}$ . The delocalization of unpaired electrons over the benzene ring may be responsible for the inertness upon H-abstraction. However, CHD has been known as a potential H donor for the excited-state species: 94% of the quenching of benzophenone triplet ( $T_1$ ) by CHD results in the H-abstraction with the rate constant of  $2.73 \times 10^8 \text{ M}^{-1} \text{ s}^{-1}$  in benzene at 300 K.<sup>17</sup> Furthermore, recent double-pulse experiments<sup>4,5</sup> have reported the fluorescence quenching processes of diphenylmethyl and 1-naphthylmethyl radicals by CHD, suggesting an excited-state H-abstraction. Hence, the quenching of a relatively long-lived fluorescence of *p*-X-benzyl (X = OCH<sub>3</sub>, F, Cl, or CN) by CHD was examined to estimate the excited-state H-abstraction reactivity. Unfortunately, the fluorescent-state benzyl radical in solution at room temperature is too short-lived ( $\leq 1$  ns) to undergo appreciable quenching by CHD.

The green  $D_1 \rightarrow D_0$  fluorescence was induced by the 308-nm XeCl laser pulse excitation of the transient *p*-methoxybenzyl formed in the 248-nm KrF laser photolysis of *p*-(chloromethyl)anisole (CMA) in hexane at room temperature. Such two-step laser-induced fluorescence (TS LIF) signals were also measured at 500 nm in the presence of high concentrations of CHD. Figure 1a shows the oscilloscope traces (A and B) for the 500-nm TS LIF with and without CHD (4360 mM). Since CHD is not transparent at the photolysis wavelength (248 nm) but transparent at the probe laser one (308 nm), the reduced TS LIF ( $D_1 \rightarrow D_0$  fluorescence) intensity may be ascribed to the filter effect of CHD for the 248-nm laser photolysis of CMA to generate *p*-methoxybenzyl. Also shown in Figure 1a is the quenched TS

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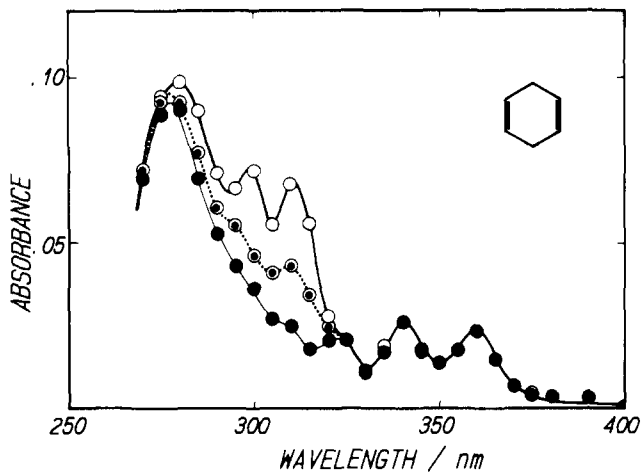
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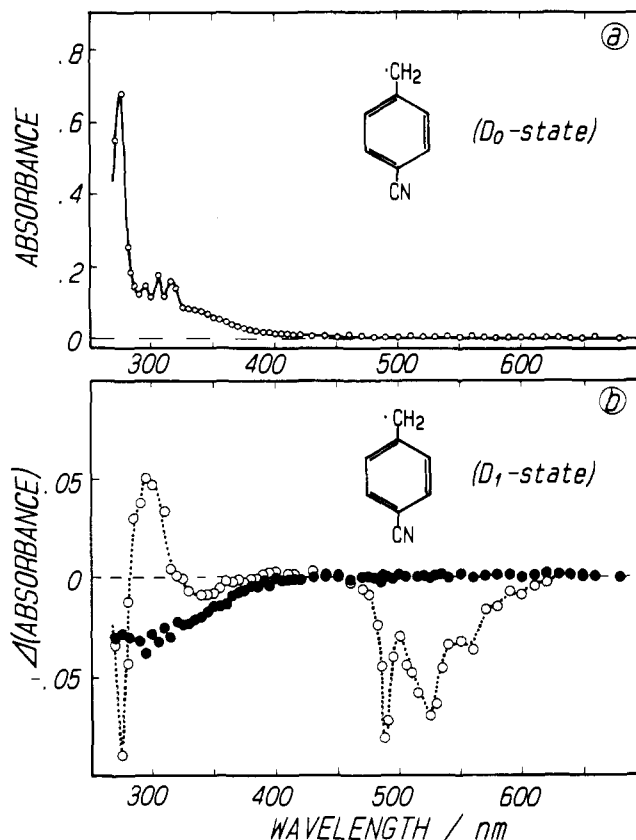


**Figure 2.** Time-resolved absorption spectra observed at 8  $\mu\text{s}$  (O), 28  $\mu\text{s}$  (⊙), and 200  $\mu\text{s}$  (●) after the 248-nm laser pulse excitation of CHD (800 mM) in degassed hexane at room temperature.

LIF signal (B'), the intensity of which is normalized at maximum with the signal (A) in the absence of CHD in contrast to the decay feature. All the TS LIF signals exhibited a single exponential decay. The fluorescence lifetimes ( $\tau$ ) were determined in the various concentrations of CHD, and the Stern-Volmer plot of  $\tau_0/\tau$  is shown in the inset of Figure 1a. The quenching rate constant ( $k_q^*$ ) was evaluated to be  $9.4 \times 10^5 \text{ M}^{-1} \text{ s}^{-1}$  from the slope of the Stern-Volmer plot. It is comparable to the rate constants ( $1.1 \times 10^6$  and  $\sim 5 \times 10^5 \text{ M}^{-1} \text{ s}^{-1}$  in cyclohexane at 300 K) reported for the fluorescence quenching of diphenylmethyl<sup>4</sup> and 1-naphthylmethyl<sup>5</sup> radicals, the unpaired electron of which can be considerably delocalized over  $\pi$ -conjugation system.

Whereas, the fluorescence of *p*-cyanobenzyl was efficiently quenched by small concentrations (3.63–14.5 mM) of CHD, as shown in Figure 1b, the  $k_q^*$  for the quenching of *p*-cyanobenzyl fluorescence was determined to be  $3.2 \times 10^9 \text{ M}^{-1} \text{ s}^{-1}$ , which is three orders of magnitude greater than that for the quenching of *p*-methoxybenzyl fluorescence. A relatively large  $k_q^*$  of  $2.0 \times 10^8 \text{ M}^{-1} \text{ s}^{-1}$  and a relatively small  $k_q^*$  of  $6.3 \times 10^6 \text{ M}^{-1} \text{ s}^{-1}$  were obtained for the fluorescence quenching of *p*-chlorobenzyl and *p*-fluorobenzyl, respectively. A marked para-substituent effect was thus confirmed for the fluorescence quenching of benzyl radical. The  $D_n \leftarrow D_1$  absorption spectra as well as the bleaching of the  $D_n \leftarrow D_0$  absorption spectra of these para-substituted benzylic would help in an interpretation of the marked substituent effect upon the reactions involved in the fluorescence quenching by CHD. Such spectra are obtainable by the two-step laser-induced-absorption measurements described in the next section.

**Two-Step Laser-Induced-Absorption Measurements for the Quenching of Fluorescent Excited-State *p*-Cyanobenzyl by 1,4-Cyclohexadiene.** Since no green emission was observed upon the 248-nm pulse excitation of CHD followed by the 308-nm pulse excitation, the quenching of the *p*-X-benzyl fluorescence by CHD could be observed as described above. However, the 248-nm pulse excitation of CHD generates the transient species subjected to the subsequent 308-nm pulse excitation. Figure 2 shows the time-resolved absorption spectra observed upon the KrF laser photolysis of CHD (800 mM) in hexane at room temperature. The residual spectrum after the decay of the absorption bands around 300 nm is almost identical with the absorption spectrum of the photolyzed solution. The unknown photoproduct is therefore responsible for the residual spectrum at 270–400 nm. There are characteristic twin peaks at 308 and 314 nm in the published UV spectrum of cyclohexadienyl radical in adamantane matrix at 298 K.<sup>18</sup> The triplet-sensitized  $T_n \leftarrow T_1$  absorption spectrum ( $\lambda_{\text{max}} = 302.5 \text{ nm}$ )<sup>19</sup> has been reported for 1,3-cyclohexadiene in



**Figure 3.** (a) Ordinary transient absorption spectrum observed at 330 ns after the 248-nm photolysis of the degassed hexane solution of  $\alpha$ -bromo-*p*-tolunitrile (0.09 mM) at room temperature. (b) Time-resolved two-step laser-induced-absorption spectra observed at 20-ns (O) and 280-ns (●) after the 308-nm laser pulse excitation delayed by 600 ns from the 248-nm laser photolysis of the solution.

benzene. The observed UV transient absorption band may be ascribed to cyclohexadienyl radical and slightly longer lived  $T_1$  state of CHD. However, the initial transient absorbance at 308 nm is less than 0.06. In order to reduce the direct photolysis of CHD as little as possible, an efficient quenching system requiring low concentration of CHD should be selected. The quenching of the excited-state *p*-cyanobenzyl by CHD was thus studied by the two-step laser-induced-absorption measurements.

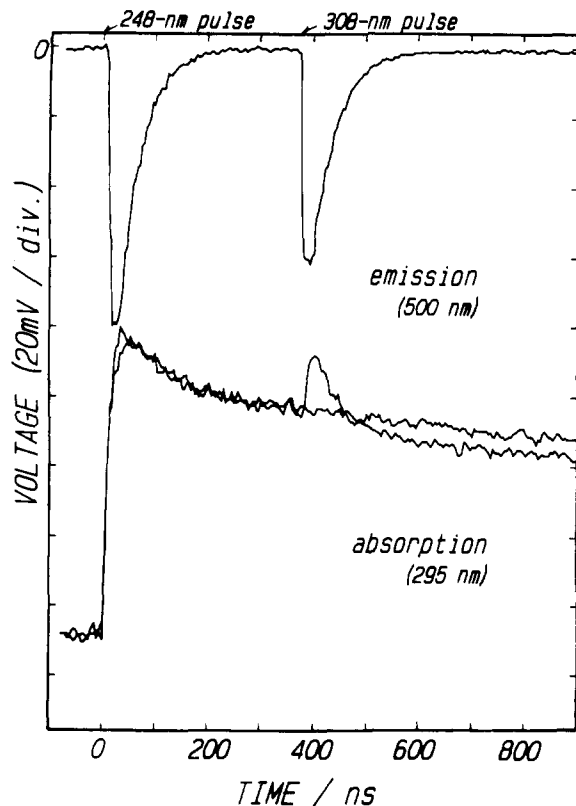
Figure 3a shows the transient absorption spectrum observed upon the 248-nm laser photolysis of  $\alpha$ -bromo-*p*-tolunitrile in degassed hexane at room temperature. The spectrum consisting of the intense band ( $\lambda_{\text{max}} = 270 \text{ nm}$ ) with weaker bands ( $\lambda_{\text{max}} = 305$  and  $315 \text{ nm}$ ) is similar to the published spectrum ( $\lambda_{\text{max}} = 280, 310, 322 \text{ nm}$ )<sup>20</sup> of *p*-cyanobenzyl formed by the reaction of  $e_{\text{aq}}^-$  with *p*-cyanobenzyl bromide followed by dehalogenation. It has been confirmed<sup>10</sup> that the  $D_1 \rightarrow D_0$  fluorescence is induced upon both the 248-nm and the 308-nm pulse excitation of *p*-methoxybenzyl formed in the 248-nm laser photolysis of *p*-(chloromethyl)anisole. Figure 4 shows the 500-nm fluorescence signals induced by the 248-nm pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile followed by the 308-nm pulse excitation. Both LIF signals exhibiting the same exponential decay of 58 ns are attributable to the  $D_1 \rightarrow D_0$  fluorescence of *p*-cyanobenzyl. It is thus reasonable to assume that the  $D_n \leftarrow D_1$  absorption signals revealed via the biphotonic processes are superimposed upon the  $D_n \leftarrow D_0$  absorption signal exhibiting a second-order decay. Such a transient absorption signal was actually observed at 295 nm in the 248-nm pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile followed by the subsequent 308-nm pulse excitation.

The decaying TS LIA signal does not join in the transient absorption signal induced by the 248-nm pulse alone, and a bleaching persists. The TS LIA spectra can be depicted by plotting

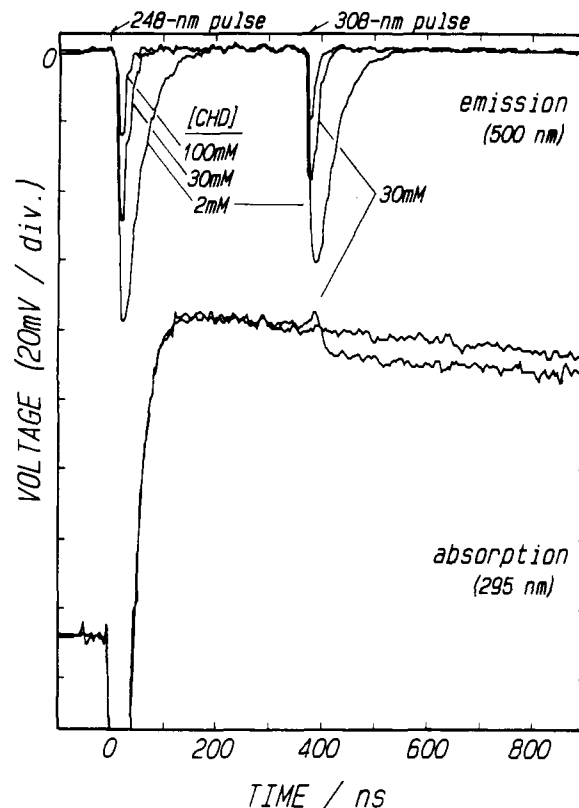
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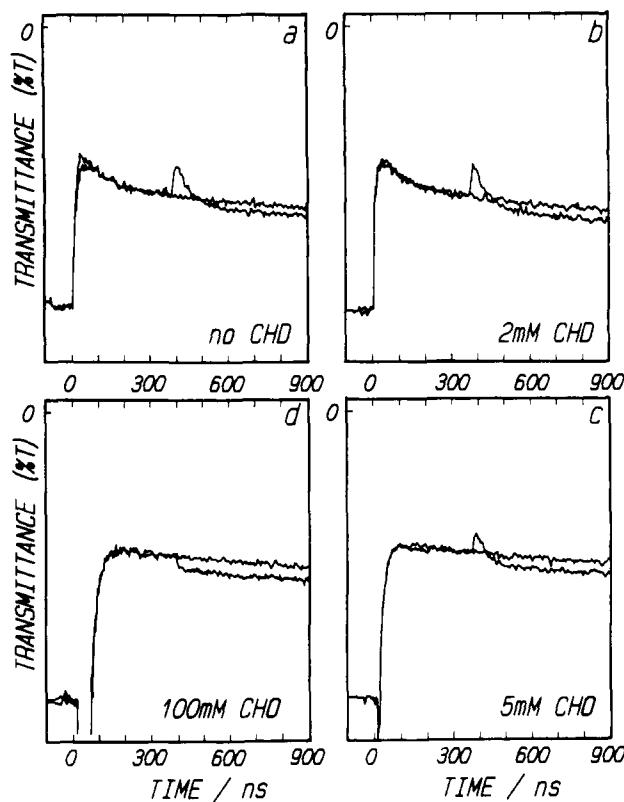
**Figure 4.** Oscillogram traces of the 295-nm transient absorption and the 500-nm fluorescence signals induced by the 248-nm laser pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile (0.09 mM) followed by the 308-nm laser pulse excitation. Also shown is an oscillogram trace of the 295-nm transient absorption induced by the 248-nm laser pulse excitation alone.



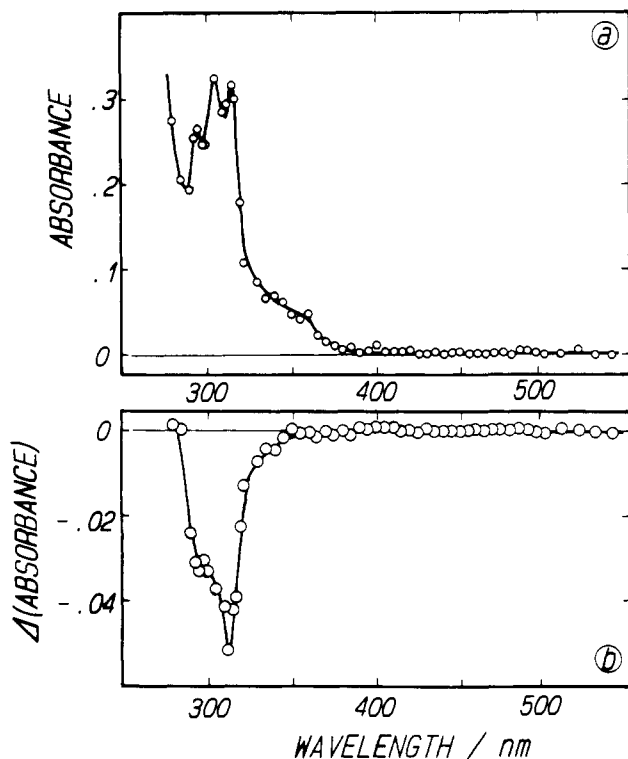
**Figure 5.** Oscillogram traces of the 295-nm transient absorption induced by the 248-nm laser pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile (0.09 mM) with and without the subsequent 308-nm laser pulse excitation. Also shown are the 500-nm LIF signals in the presence of CHD (2, 30, 100 mM): first laser pulse, 248 nm; second laser pulse, 308 nm.

log  $[I_1(t)/I_2(t)]$  against wavelength. Here,  $I_1(t)$  and  $I_2(t)$  stand for the monitoring light intensity for the ordinary (one-step) LIA and TS LIA signals, respectively. Figure 3b shows the TS LIA spectra at 20 and 280 ns after the 308-nm pulse excitation delayed by ca. 600 ns from the 248-nm pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile. The spectrum at 20-ns delay was actually depicted by plotting log  $(I_1/I_2(t))$  against wavelength. Here, the monitoring light intensity immediately before the 308-nm pulse excitation is designated as  $I_1$ . The  $D_1 \rightarrow D_0$  fluorescence is responsible for the apparent negative absorption band (450–650 nm), and the UV TS LIA spectrum may be primarily recognized as the difference spectrum between the  $D_n \leftarrow D_1$  and the  $D_n \leftarrow D_0$  absorption spectra. The positive absorption band ( $\lambda_{\max} = 295$  nm) is thus regarded as the  $D_n \leftarrow D_1$  absorption spectrum of *p*-cyanobenzyl, and the sharp negative band at 270 nm is due to the depletion of the intense  $D_n \leftarrow D_0$  absorption peak.

The TS LIA spectrum was also taken at 280 ns after the 308-nm pulse excitation. At that time, the  $D_1 \rightarrow D_0$  fluorescence has completely faded out. The negative absorption spectrum is similar to the inversion of the 330–400-nm broad absorption band (Figure 3a) due to unknown species formed in the 248-nm laser photolysis of  $\alpha$ -bromo-*p*-tolunitrile. No significant bleaching of the  $D_n \leftarrow D_0$  absorption is seen in this spectrum. Figure 5 shows the 500-nm  $D_1 \rightarrow D_0$  fluorescence signals of *p*-cyanobenzyl induced by the 248-nm pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile followed by the 308-nm pulse excitation in the presence of CHD (2, 30, and 100 mM). The 295-nm transient absorption in the presence of 30 mM of CHD is also shown in Figure 5. Furthermore, Figure 6 shows the 295-nm TS LIA as well as the ordinary (one-step) LIA in the absence of CHD (a) and in the presence of CHD (b–d: 2, 5, 100 mM). The quenching of the excited-state *p*-cyanobenzyl is evident from the fact that the TS LIA ( $D_n \leftarrow D_1$  absorption) decays faster with increasing concentration of CHD. In the considerable concentrations of CHD, the  $D_n \leftarrow D_1$  absorption signal is hidden by the emission of CHD, as shown in Figure 6c,d.



**Figure 6.** Oscillogram traces of the 295-nm transient absorption induced by the 248-nm laser pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile (0.09 mM) with and without the subsequent 308-nm laser pulse excitation in the absence of CHD (a) and in the presence of CHD: (b) 2 mM, (c) 5 mM, (d) 100 mM.



**Figure 7.** (a) Ordinary transient absorption spectrum observed at 300 ns after the 248-nm laser photolysis of  $\alpha$ -bromo-*p*-tolunitrile (0.09 mM) in the presence of 100 mM of CHD in hexane at room temperature. (b) Two-step laser-induced-absorption spectrum observed at 140 ns after the 308-nm laser pulse excitation delayed by 370 ns from the 248-nm laser photolysis of the solution.

Figure 7a shows the transient absorption spectrum observed at 300 ns after the 248-nm pulse excitation of the solution of  $\alpha$ -bromo-*p*-tolunitrile and CHD (100 mM).<sup>21</sup> The spectrum is similar to that (Figure 3a) in the absence of CHD. Three peaks around 300 nm are ascribed to the  $D_n(2A_2) \leftarrow D_0(1B_2)$  absorption spectrum of *p*-cyanobenzyl, and the weak and broad absorption band around 350 nm may be due to unknown species,<sup>22</sup> as already described earlier. The H-abstraction of *p*-cyanobenzyl ( $D_1$ ) from a CHD is assumed to generate a *p*-tolunitrile and a cyclohexadienyl radical. Almost the same absorption spectra of *p*-tolunitrile and  $\alpha$ -bromo-*p*-tolunitrile prevent us from detecting the formation of *p*-tolunitrile in the TS-LIA measurements. The transient visible absorption spectrum ( $\lambda_{\max} = 560$  nm)<sup>23</sup> of cyclohexadienyl radical is not detectable owing to its poor molar extinction coefficient ( $<100$  M<sup>-1</sup> cm<sup>-1</sup>).<sup>4</sup> However, the aforementioned UV spectrum of cyclohexadienyl is much more intense than the visible band.

The molar extinction coefficient of cyclohexadienyl has been reported to be  $1.21 \times 10^4$  M<sup>-1</sup> cm<sup>-1</sup> at 318 nm,<sup>24</sup> while that of *p*-cyanobenzyl is ca.  $5.0 \times 10^3$  M<sup>-1</sup> cm<sup>-1</sup> at 322 nm.<sup>20</sup> If the H-abstraction of the excited-state ( $D_1$ ) *p*-cyanobenzyl from CHD to yield cyclohexadienyl is predominant, the TS-LIA spectrum

(21) The 248-nm laser photolysis of CHD (800 mM) gave rise to weak transient absorption spectra ascribable to the lowest triplet state and cyclohexadienyl radical, as described in the text. The 248-nm laser pulse excitation of CHD is much less for the solution containing both  $\alpha$ -bromo-*p*-tolunitrile (0.09 mM) and CHD (100 mM) than for the solution of CHD (800 mM) alone.

(22) Methylated benzyl cations are known to be formed as a minority product in the UV steady-state illumination of methylated benzenes in rigid glass at 77 K. On the other hand, long-lived and broad spectra around 300 nm were confirmed in the transient absorption spectra of all *p*-X-benzyl halides examined. It is highly probable that the long-lived transient absorption is due to the *p*-X-benzyl cation formed by a consecutive and/or a simultaneous two 248-nm photon absorption of *p*-X-benzyl halide within the duration of KrF laser pulse.

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**Table I.** Rate Constants<sup>a</sup> for the Fluorescence Quenching of Para-Substituted Benzyl Radicals by CHD in Hexane at 300 K

X ( <i>p</i> -X-benzyl)	$k_q^*/M^{-1} s^{-1}$	$D_1$ state
-CN	$3.2 \times 10^9$	2B <sub>2</sub>
-Cl	$2.0 \times 10^8$	1A <sub>2</sub>
-F	$6.3 \times 10^6$	1A <sub>2</sub>
-OCH <sub>3</sub>	$9.4 \times 10^5$	1A <sub>2</sub>

<sup>a</sup>Standard deviation ca. 15%.

around 320 nm should be fairly positive regardless to the negative absorption reflecting the bleaching of the weak spectrum due to unknown species. The persistent bleaching spectrum (Figure 7b) was actually obtained at 140 ns after the 308-nm pulse excitation delayed by 370 ns from the 248-nm pulse excitation of  $\alpha$ -bromo-*p*-tolunitrile. It is considerably different from the bleaching spectrum (Figure 3b), reflecting the depletion of the unknown photoproduct in the absence of CHD, and it rather resembles the inversion of the  $D_n(2A_2) \leftarrow D_0(1B_2)$  absorption spectrum. The lack of the bleaching peak corresponding to the 305-nm peak of the  $D_n(2A_2) \leftarrow D_0(1B_2)$  absorption spectrum of *p*-cyanobenzyl implies the formation of cyclohexadienyl exhibiting a considerable absorption around 300 nm. From these facts, it is demonstrated that some fractions of the quenching actions of the excited-state *p*-cyanobenzyl by CHD result in the H-abstraction reaction to yield cyclohexadienyl.

The photoselection method applied for benzyl and *p*-methylbenzyl in rigid glass at 77 K has indicated the assignment of 1A<sub>2</sub> for their emitting state.<sup>25</sup> The energy levels of the close-lying lowest excited doublet states (1A<sub>2</sub> and 2B<sub>2</sub>) of benzyl may be perturbed by some substituents. Contrary to only a minor perturbation of the methyl group or chlorine atom at the para position of benzyl,<sup>12</sup> both *p*-methoxy<sup>12</sup> and *p*-cyano<sup>20</sup> groups make a strong perturbation to the electronic levels of benzyl system. The  $D_1 \rightarrow D_0$  fluorescence spectra of *p*-methoxybenzyl and *p*-cyanobenzyl in rigid glass at 77 K exhibited the 0,0-band at 463 nm and 461 nm, respectively.<sup>26</sup> Nevertheless, we actually observed the considerably different  $D_n \leftarrow D_1$  absorption spectra exhibiting the peak at 340 nm<sup>10</sup> for *p*-methoxybenzyl and at 295 nm for *p*-cyanobenzyl. Assuming that the allowed  $D_n$  state is severely dependent upon the nature of the lowest excited doublet state ( $D_1$ ), the different  $D_1$  states may be responsible for their quite different  $D_n \leftarrow D_1$  absorption spectra. Very recently, Hiratsuka et al.<sup>27-29</sup> have determined the polarization degree of the fluorescence emission and excitation spectra of *p*-cyanobenzyl, *p*-carboxybenzyl, *p*-fluorobenzyl, and *p*-methoxybenzyl as well as benzyl itself stabilized in the polyvinyl alcohol film at 77 K. They assigned the emitting state of *p*-cyanobenzyl and *p*-carboxybenzyl to 2B<sub>2</sub> contrary to the 1A<sub>2</sub> assignment for other benzyl radicals, and they suggested that the strong  $\pi$ -electron-attracting group at the para position of benzyl causes an inversion of the close-lying 1A<sub>2</sub> and 2B<sub>2</sub> states. These facts are very consistent with the significantly different  $D_n \leftarrow D_1$  absorption spectra of *p*-cyanobenzyl and *p*-methoxybenzyl described above.

A marked substituent effect was demonstrated for the quenching of the *p*-X-benzyl fluorescence and for the assignment of the lowest excited doublet state, as summarized in Table I. Although all the quenching actions of the excited-state *p*-X-benzyl by CHD do not lead to the H-abstraction from CHD, the excited-state H-abstraction reactivity is the primary factor for the fluorescence quenching by CHD. The marked substituent effect upon the quenching by CHD is therefore ascribed to the difference in the unpaired electron density, and it is assumed that the unpaired

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electron density at the benzyl position of  $2B_2$  is considerably higher than that of  $1A_2$  state. Significant mixing of the close-lying  $2B_2$  ( $D_2$ ) and  $1A_2$  ( $D_1$ ) might be responsible for a relatively high rate in the quenching of the *p*-chlorobenzyl fluorescence by CHD. From the analysis of the vibrational structure in the fluorescence excitation spectrum of *p*-fluorobenzyl in a supersonic jet, Fukushima and Obi<sup>30</sup> have demonstrated that the energy gap between  $1A_2$  and  $2B_2$  states is spread out by the introduction of fluorine atom at the para position of benzyl. Poor mixing of the two states might be therefore responsible for a low rate in the quenching

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of *p*-fluorobenzyl fluorescence by CHD. The predominant character of  $1A_2$  is implied for the  $D_1$  state of *p*-methoxybenzyl by a very low rate of  $9.4 \times 10^5 \text{ M}^{-1} \text{ s}^{-1}$  in the fluorescence quenching.

**Acknowledgment.** This work was supported by the Grant-in-Aid (No. 63104007) for Scientific Research from the Ministry of Education, Science and Culture of Japan. The authors are grateful to Dr. H. Hiratsuka for his courtesy to give us the manuscript prior to publication.

**Registry No.** CHD, 628-41-1; *p*-cyanobenzyl radical, 4939-73-5; *p*-chlorobenzyl radical, 3327-51-3; *p*-fluorobenzyl radical, 2194-09-4; *p*-methoxybenzyl radical, 3494-45-9.

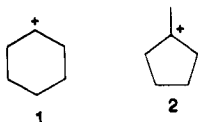
## Evaluation of the Arrhenius Parameters for the Thermal Isomerization of a Gaseous Ion. The Temperature Dependence of the Cyclohexylium Ion Unimolecular Rearrangement

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Received October 20, 1988

**Abstract:** Gas-phase isomerization of cyclohexylium ion (**1**), obtained via protonation of cyclohexene and of bicyclo[3.1.0]hexane and via hydride-ion abstraction from cyclohexane by  $C_nH_5^+$  ( $n = 1, 2$ ) and by  $s\text{-C}_3\text{H}_7^+$  ions, has been investigated by a radiolytic technique in  $\text{CH}_4$  and  $\text{C}_3\text{H}_8$  in the pressure range from 106 to 1480 Torr. Temperature-dependence studies in the temperature interval from 310 to 353 K have provided the activation parameters of the thermal rearrangement of **1** into 1-methylcyclopentylum ion (**2**) in  $\text{CH}_4$  at 750 Torr and in  $\text{C}_3\text{H}_8$  at 1480 Torr. Least-squares analysis of the two mutually consistent sets of data leads to an activation energy of  $7.4 \pm 1 \text{ kcal mol}^{-1}$  and to a preexponential factor of  $10^{12 \pm 1.3} \text{ s}^{-1}$ . These values represent the first estimates of the activation parameters of a unimolecular isomerization of a gaseous ion thermally equilibrated with the bath gas, i.e., under conditions which allow meaningful comparison with solution kinetics. The results support previous evidence for the existence of gaseous **1**, providing a quantitative evaluation of the energy barrier for its rearrangement into **2**. The more extensive formation of **2** from the protonation of bicyclo[3.1.0]hexane than that of cyclohexene is consistent with the long-postulated protonated-cyclopropane route in the branching rearrangement of cycloalkylium ions.

Cyclohexylium cation (**1**), well-characterized as an ionic intermediate in solution,<sup>1</sup> has long escaped detection in the gas phase, owing to its fast rearrangement into the 1-methylcyclopentylum isomer (**2**), occurring in a shorter time than required for analysis



by structurally diagnostic mass spectrometric techniques, e.g., collisional activation (CA) spectrometry.<sup>2</sup> In previous papers<sup>3,4</sup> we have presented evidence for the existence of **1** in the gas phase, with a lifetime exceeding  $10^{-8}$ – $10^{-7}$  s, based on the results obtained with a radiolytic<sup>5</sup> and a nuclear-decay technique,<sup>6</sup> which allow

structural and kinetic studies of gaseous ions in a wide pressure range. While demonstrating that **1** is indeed a legitimate ionic intermediate in the gas phase, characterized by a local minimum on the  $C_6H_{11}^+$  energy surface,<sup>7</sup> our results confirmed its facile rearrangement into structure **2**, more stable by 11 kcal mol<sup>-1</sup>,<sup>8</sup> the activation energy of the process being roughly estimated as  $\leq 10 \text{ kcal mol}^{-1}$ .<sup>4</sup>

We report now a more detailed investigation of the gas-phase isomerization of **1**, including an attempt at estimating its empirical activation parameters by temperature-dependence experiments in the pressure range from 106 to 1480 Torr.

It should be emphasized that meaningful evaluation of activation parameters comparable to those ordinarily measured in solution requires that the ions undergoing isomerization be thermally equilibrated with a bath gas of defined temperature, and that their rearrangement can be monitored by sufficiently fast techniques. These conditions can hardly be achieved by current mass spectrometric approaches, where the ionization process imparts excess internal energy to the charged species, whose collisional thermalization is generally inefficient at the low pressures attainable

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